

# RBI Quantum Hackathon Workbench - this time on Learning Parity with Noise from Machine Learning Viewpoint

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## Revision History

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- 06/10/2019/Tom - LPN introduction, Nature-style experiment elaboration
- 07/10/2019/Tom - Generalised ancilla initialisation leading to a better algorithm
- 08/10/2019/Tom - Improved readability by detailing key formulas derivations
- 17/10/2019/Tom - Epsilon ancilla qubit incorporated

# Notation Notes

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- $\oplus$  denotes (vector) addition modulo 2
- $\odot$  denotes (vector) multiplication modulo 2
  - for binary vectors,  $\odot$  is a standard dot product modulo 2
  - when clear from the context, we use simply  $+$  and  $\oplus$ , or  $\cdot$  and  $\odot$  interchangeably

# Symmetrisation Intermezzo

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- We will often work with output of boolean functions like  $f(x)$ , where  $x$  is a binary vector
- This output is in general either 0 or 1
- When designing quantum algorithms, we need to incorporate  $f(x)$  into superposition coefficients in a concise way to see the effect of quantum operators
- For this, it is useful to transform  $f(x)$  so to make its result either -1 or 1, or to incorporate the effect of  $f(x)$  through a power of (-1)

# Symmetrisation We Use

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Let  $f(x) \in \{0,1\}$ .

Then:

$$2f(x) - 1 \in \{-1,1\}, \quad -1 \text{ iff } f(x) = 0$$

$$(-1)^{f(x)} \in \{-1,1\}, \quad -1 \text{ iff } f(x) = 1$$

In particular:

$$(-1)^{f(x)} = 1 - 2f(x)$$

$$1 + (-1)^{f(x)} = 2(1 - f(x))$$

$$1 - (-1)^{f(x)} = 2f(x)$$

Also note the Hadamard transform of such a boolean  $f(x)$

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$$|f(x)\rangle \mapsto \frac{(-1)^{0 \cdot f(x)} |0\rangle + (-1)^{1 \cdot f(x)} |1\rangle}{\sqrt{2}}$$

$$= \frac{|0\rangle + (-1)^{f(x)} |1\rangle}{\sqrt{2}}$$

# Learning Parity with Noise

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The *search* version of the learning parity with noise problem with parameters  $\ell \in \mathbb{N}$  (the length of the secret),  $\tau \in \mathbb{R}$  where  $0 < \tau < 0.5$  (the noise rate) and  $q \in \mathbb{N}$  (the numbers of samples) asks to find a fixed random  $\ell$  bit secret  $\mathbf{s} \in \mathbb{Z}_2^\ell$  from  $q$  samples of the form  $\mathbf{a}, \langle \mathbf{a}, \mathbf{s} \rangle \oplus e$  where  $\mathbf{a} \in \mathbb{Z}_2^\ell$  is random and  $e \in \mathbb{Z}_2$  has Bernoulli distribution with parameter  $\tau$  (we denote this distribution with  $\text{Ber}_\tau$ ), i.e.  $\Pr[e = 1] = \tau$ . The *decisional* LPN problem is defined similarly, except that we require that one cannot even distinguish noisy inner products from random.

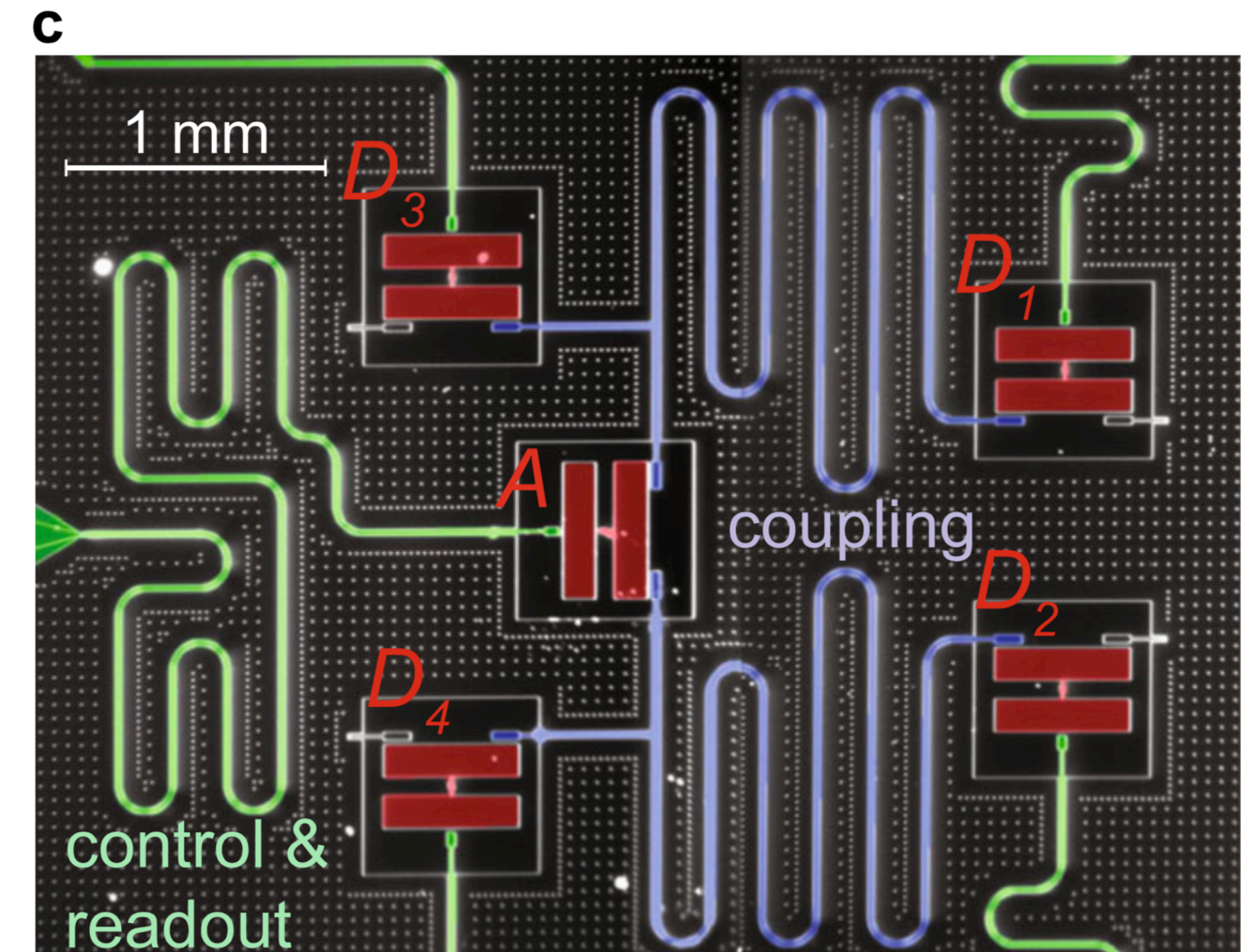
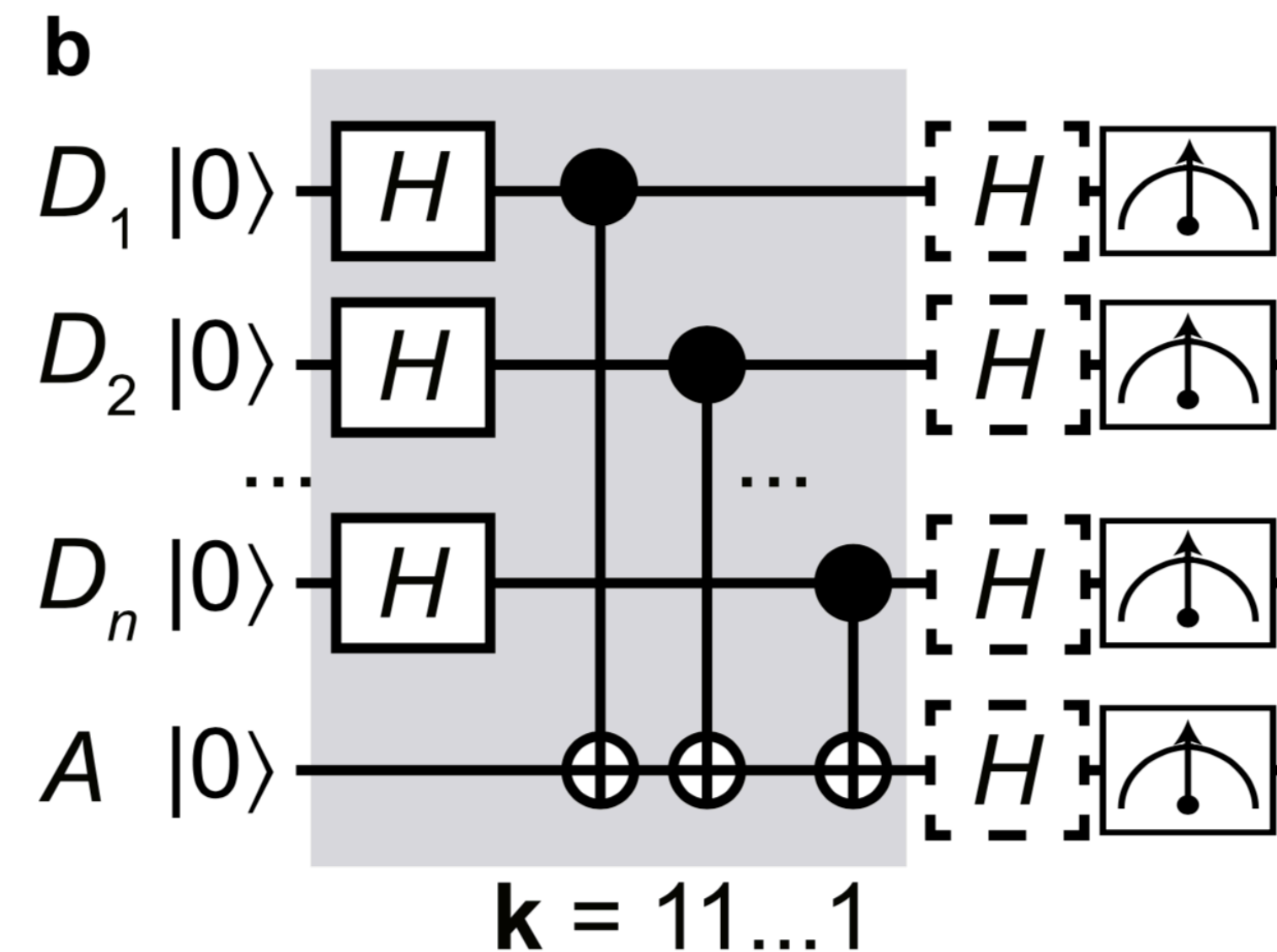
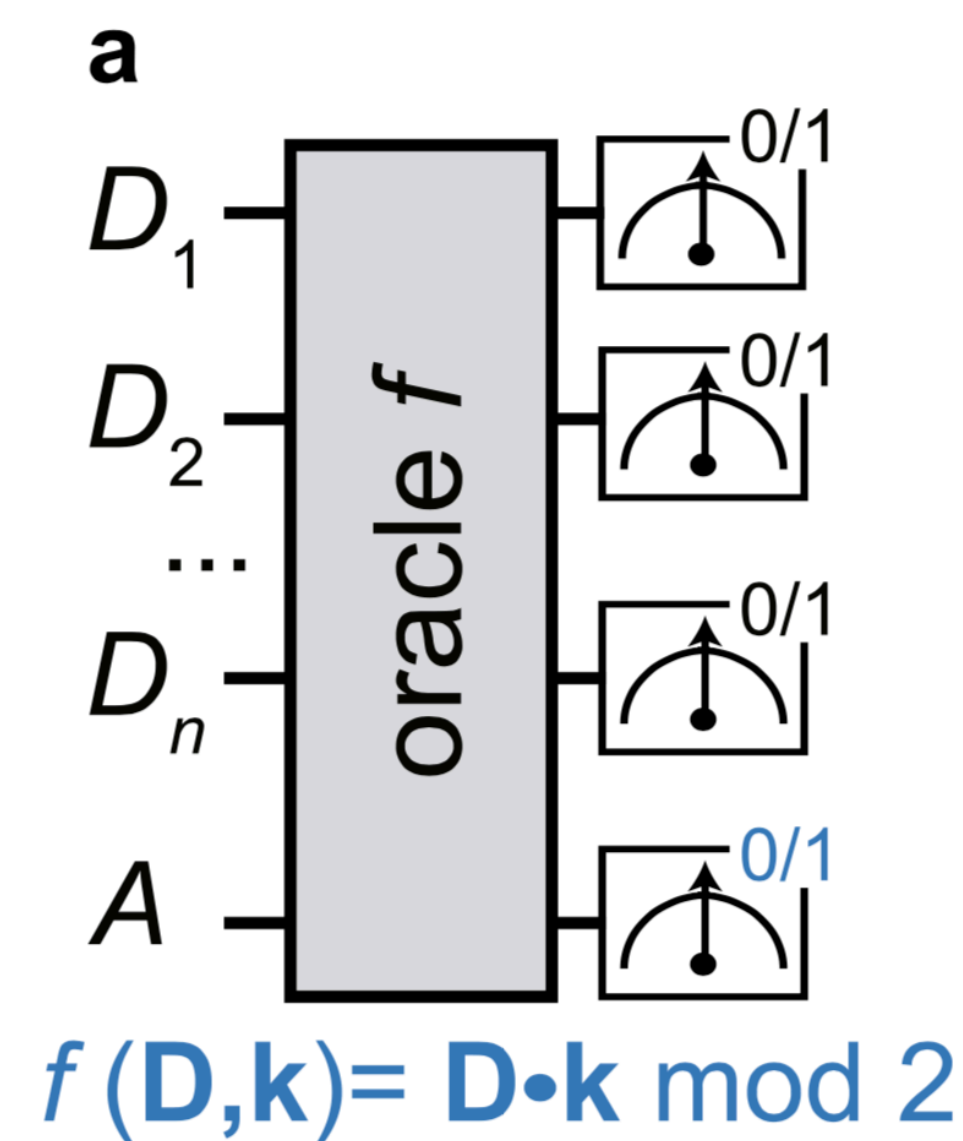
# Broad Impact of LPN

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- We have chosen the Learning Parity with Noise problem for this hackathon, since it has considerable impact on both
  - machine learning techniques
  - post-quantum cryptography and cryptanalysis



# Starting Experiment Described in Nature Partner Journal on QI (4/17)



ARTICLE **OPEN**

## Demonstration of quantum advantage in machine learning

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# Our Starting Position

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- We consider the quantum approach to LPN solving, so we always assume the final Hadamard gates are on
  - omitting output Hadamard(s) was to simulate classical LPN conditions with the same QPU (*Quantum Processing Unit*) core setup
- In the original, the LPN noise is intrinsic, generated by QPU inherently
  - we stick more with classical LPN, so we explicitly use the additive error factor
  - it is generated independently for each oracle-operator invocation

# Our Goals

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- Improve the efficiency of the original algorithm by showing its direct connection with Bernstein-Vazirani algorithm we elaborated in Vienna in May this year (so btw., there is an ongoing competence extension and application)
  - it follows from a generalisation of the ancilla qubit initialisation
  - actually, it is a bit surprising the former authors did not note this connection already
- Verify the theoretical construction practically, even with a higher number of data qubits (originally, they used two)
  - we want to show this approach practically halves the number of LPN oracle invocations (respectively the number of QPU runs); as we are talking about practical machine learning algorithms, this can be significant
- Discuss the quantum advantage
  - we should be able to solve even the worst-case LPN instances that are unsolvable classically
  - i.e. for  $\tau = 1/2$

We investigate the original (“Nature-style”) initialisation and approach

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$$|\psi_1\rangle = \frac{1}{\sqrt{2^n}} \sum_{x \in F_2^n} |x\rangle \otimes |0\rangle$$

note we have dropped the bar vector notation used before



We apply the LPN oracle operator for  $f_k(x)$

$k$  is the hidden number,  $\varepsilon$  is the binary error factor (the noise)

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$$|x\rangle \otimes |0\rangle \mapsto |x\rangle \otimes |f_k(x) \oplus 0\rangle, \text{ where } f_k(x) = k \odot x \oplus \varepsilon$$

$$|\psi_{2,k}\rangle = \frac{1}{\sqrt{2^n}} \sum_{x \in F_2^n} |x\rangle \otimes \left( \frac{1 + (-1)^{f_k(x)}}{2} |0\rangle + \frac{1 - (-1)^{f_k(x)}}{2} |1\rangle \right)$$



phase kickback effect **not obvious**, now

Then we apply the final Hadamard transform(s) (cf. original scheme)

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$$|\psi_{3,k}\rangle = \frac{1}{2^n} \sum_{y \in F_2^n} \sum_{x \in F_2^n} (-1)^{x \odot y} |y\rangle \otimes \left[ \frac{1 + (-1)^{f_k(x)}}{2\sqrt{2}} (|0\rangle + |1\rangle) + \frac{1 - (-1)^{f_k(x)}}{2\sqrt{2}} (|0\rangle - |1\rangle) \right]$$

$$= \frac{1}{2^n} \sum_{y \in F_2^n} \sum_{x \in F_2^n} (-1)^{x \odot y} |y\rangle \otimes \left( \frac{|0\rangle + (-1)^{f_k(x)} |1\rangle}{\sqrt{2}} \right), \text{ for } f_k(x) = k \odot x \oplus \varepsilon$$

$$= \frac{1}{2^n \sqrt{2}} \sum_{y \in F_2^n} \sum_{x \in F_2^n} (-1)^{x \odot y} |y\rangle |0\rangle + \frac{1}{2^n \sqrt{2}} \sum_{y \in F_2^n} \sum_{x \in F_2^n} (-1)^\varepsilon (-1)^{x \odot (y \oplus k)} |y\rangle |1\rangle$$

$$= \frac{1}{\sqrt{2}} \underbrace{|00\dots 0\rangle}_{\text{garbage}} |0\rangle + \frac{(-1)^\varepsilon}{\sqrt{2}} \underbrace{|k\rangle}_{\text{direct secret bits}} |1\rangle$$

garbage

direct secret bits

# Notes on the Original Approach

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- We have 50% chance to measure  $|0\rangle$  and  $|1\rangle$  on the ancilla qubit, respectively
  - measuring  $|0\rangle$  brings no further information in data qubits
  - measuring  $|1\rangle$  reveals the hidden number (secret)  $k$  in data qubits
- In case of the positive answer, the result is totally insensitive on the  $\varepsilon$  noise
  - this is certainly a good point (though not so much addressed before)
- We, however, waste around 50% of QPU runs
  - we shall try to do this better, now

# In Search for the Generalised Ancilla Initialisation (GAI)

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- Note that the original “Nature-style” approach works practically the same way, regardless whether the ancilla is initialised as  $|0\rangle$  or  $|1\rangle$ 
  - this suggests that both of the pure eigenstates are equally good or bad
  - how about to try their superposition?
  - heuristically, we try an equal superposition with a variable relative phase



# The Generalised Ancilla Initialisation (GAI)

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$$|\psi_1\rangle = \frac{1}{\sqrt{2^n}} \sum_{x \in F_2^n} |x\rangle \otimes \left( \frac{|0\rangle + e^{i\varphi} |1\rangle}{\sqrt{2}} \right)$$

Again, we apply the LPN oracle operator for  $f_k(x)$

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$$|\psi_{2,k}\rangle = \frac{1}{\sqrt{2^n}} \sum_{x \in F_2^n} |x\rangle \otimes \left( \frac{|f_k(x)\rangle + e^{i\varphi} |f_k(x) \oplus 1\rangle}{\sqrt{2}} \right)$$



phase kickback can be shown for e.g.  $\varphi = \pi$

Now, we use the final Hadamard transform(s) to see how it works

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$$\begin{aligned}
 |\psi_{3,k}\rangle &= \frac{1}{2^n} \sum_{y \in F_2^n} \sum_{x \in F_2^n} (-1)^{x \odot y} |y\rangle \otimes \left( \frac{|0\rangle + (-1)^{f_k(x)} |1\rangle}{2} + e^{i\varphi} \frac{|0\rangle + (-1)^{f_k(x) \oplus 1} |1\rangle}{2} \right) \\
 &= \frac{1}{2^n} \sum_{y \in F_2^n} \sum_{x \in F_2^n} (-1)^{x \odot y} |y\rangle \otimes \left( \frac{(1+e^{i\varphi})|0\rangle + (1-e^{i\varphi})(-1)^{f_k(x)} |1\rangle}{2} \right), \text{ note } (-1)^{f_k(x) \oplus 1} = -(-1)^{f_k(x)} \\
 &= \frac{1+e^{i\varphi}}{2^{n+1}} \sum_{y \in F_2^n} \sum_{x \in F_2^n} (-1)^{x \odot y} |y\rangle |0\rangle + \frac{1-e^{i\varphi}}{2^{n+1}} \sum_{y \in F_2^n} \sum_{x \in F_2^n} (-1)^\varepsilon (-1)^{x \odot (y \oplus k)} |y\rangle |1\rangle, \text{ note } f_k(x) = k \odot x \oplus \varepsilon \\
 &= \frac{1+e^{i\varphi}}{2} \underbrace{|00\dots 0\rangle}_{\text{garbage}} |0\rangle + \frac{(-1)^\varepsilon (1-e^{i\varphi})}{2} \underbrace{|k\rangle}_{\text{direct secret bits}} |1\rangle
 \end{aligned}$$

# Minimising the Garbage Probability

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- Using GAI, we can manipulate the probability of the garbage  $|00\dots 0\rangle|0\rangle$  state observation
  - this way, we can minimise the waste of QPU runs, so to increase effectiveness of our LPN quantum solver
- The garbage probability is equal to zero if

$$1 + e^{i\varphi} = 0 \Leftrightarrow \varphi = (2b + 1)\pi, b \in \mathbb{Z}$$

# Welcome back, please, Mr. Bernstein and Mr. Vazirani

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- With  $\varphi = \pi$ , we get exactly the Bernstein-Vazirani algorithm (BVA) again
  - as reformulated by Cleve et al. in “*Quantum Algorithms Revisited*”, 1997
  - apparently, this algorithm is quite powerful for both machine learning and cryptology
  - **we have shown that BVA is a general extension of the Nature-style approach and the most efficient way to solve the LPN studied here**

# Experimental Implementation of the $\varepsilon$ -error

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- To fully implement the LPN oracle, we would need to be able to alter its quantum operator for each and every QPU run
  - hard to do with the actual Qiskit platform
- We decided to do an equivalent implementation based on an extra error-driving qubit
  - the error is still interpreted classically, but it is inserted in a quantum way
  - adding the  $|q_\varepsilon\rangle$  ancilla qubit tweaks the state after the “error-free” LPN operator using a CNOT entanglement to a superposition of error-free and erroneous substates
  - our LPN solver solves both instances in parallel; finally revealing only one, depending on the error probability distribution
  - by measuring the epsilon ancilla, we can get a kind of “debug” information for further statistical processing

# Epsilon Ancilla Qubit Before Entangling with the LPN Output (General Distribution)

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$$|\psi_{2,k}\rangle = \frac{1}{\sqrt{2^n}} \sum_{x \in F_2^n} |x\rangle \otimes \left( \frac{|f_k(x)\rangle + e^{i\varphi} |f_k(x) \oplus 1\rangle}{\sqrt{2}} \right) \otimes \left( \cos \frac{\theta}{2} |0_\varepsilon\rangle + \sin \frac{\theta}{2} e^{i\beta} |1_\varepsilon\rangle \right)$$

epsilon ancilla added; state is tweaked via tensor product

— here, the  $f_k$  is just the inner product, without the epsilon error

# CNOT Entangling the Epsilon Error with the LPN Output (General Error Distribution)

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$$\begin{aligned}
 |\psi^{(\varepsilon)}_{2,k}\rangle &= \frac{\cos \frac{\theta}{2}}{\sqrt{2^n}} \sum_{x \in F_2^n} |x\rangle \otimes \left( \frac{|f_{k,\varepsilon=0}(x)\rangle + e^{i\varphi} |f_{k,\varepsilon=0}(x) \oplus 1\rangle}{\sqrt{2}} \right) \otimes |0_\varepsilon\rangle \\
 &\quad \text{— error-free branch} \\
 &+ \frac{\sin \frac{\theta}{2} e^{i\beta}}{\sqrt{2^n}} \sum_{x \in F_2^n} |x\rangle \otimes \left( \frac{|f_{k,\varepsilon=1}(x)\rangle + e^{i\varphi} |f_{k,\varepsilon=1}(x) \oplus 1\rangle}{\sqrt{2}} \right) \otimes |1_\varepsilon\rangle \\
 &\quad \text{— erroneous branch}
 \end{aligned}$$

— here, the  $f_{k,\varepsilon}$  is the inner product with the explicit epsilon error



# What Follows

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- Standard finalisation via  $|\Psi_3\rangle$ 
  - both error-free and erroneous branches are solved in parallel
  - the measurement finally reveals either the branch for  $\varepsilon=0$  or  $\varepsilon=1$ , respectively
  - by observing the epsilon ancilla qubit, we can get further statistical discrimination to verify our solver works for both situations equally well